Environmental Fluid Dynamics: Lecture 10

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Overview

1 Atmospheric Dynamics: Basic Equations Conservation of Momentum, continued



Atmospheric Dynamics:

Conservation of Momentum, continued

- Apparent forces (centrifugal, Coriolis) will automatically appear when we write Newton's 2nd Law in a form appropriate for a rotating reference frame.
- Consider a Cartesian coordinate system x', y', and z' rotating with angular velocity $\overrightarrow{\Omega}$.
- Let \hat{i}' , \hat{j}' , \hat{k}' be unit vectors in the x', y', and z' directions, respectively (these vectors rotate with the coordinate system).



- Consider some arbitrary vector \overrightarrow{A} that is observed in an inertial reference frame.
- Even though \overrightarrow{A} is observed in a non-rotating reference frame, it can be decomposed into components in a rotating coordinate system:

$$\vec{A} = A_x'\hat{i}' + A_y'\hat{j}' + A_z'\hat{k}'$$

where

$$A'_{x} = \hat{i}' \cdot \vec{A}$$

$$A'_{y} = \hat{j}' \cdot \vec{A}$$

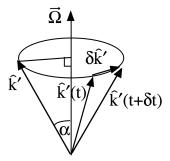
$$A'_{z} = \hat{k}' \cdot \vec{A}$$



• The total derivative of \overrightarrow{A} as observed in the inertial reference frame is (where subscript a means absolute, or inertial):

$$\begin{split} \frac{D_a \overrightarrow{A}}{Dt} &= \frac{D_a}{Dt} \left(A_x' \hat{i}' + A_y' \hat{j}' + A_z' \hat{k}' \right) \\ &= \underbrace{\hat{i}' \frac{D_a A_x'}{Dt} + \hat{j}' \frac{D_a A_y'}{Dt} + \hat{k}' \frac{D_a A_z'}{Dt}}_{D\overrightarrow{A}/Dt \text{ observed in rotating ref frame}} \\ &+ A_x' \frac{D_a \hat{i}'}{Dt} + A_y' \frac{D_a \hat{j}'}{Dt} + A_z' \frac{D_a \hat{k}'}{Dt} \\ &= \underbrace{D\overrightarrow{A}}_{Dt} + \underbrace{A_x' \frac{D_a \hat{i}'}{Dt} + A_y' \frac{D_a \hat{j}'}{Dt} + A_z' \frac{D_a \hat{k}'}{Dt}}_{\text{What are these terms?}} \end{split}$$

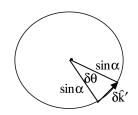
• $D_a \hat{k}'/Dt$ is the rate of change of the unit vector \hat{k}' as observed in in a inertial reference frame:



- tip of $\hat{k}'(t)$ traces out a circle, α is the angle between $\overrightarrow{\Omega}$ and $\hat{k}'(t)$, and the radius of the circle $=|\hat{k}'|\sin\alpha=\sin\alpha$.
- $\delta \hat{k}' \equiv \hat{k}'(t+\delta t) \hat{k}'(t)$ is the change in \hat{k}' over a small time interval δt .



• Let's look at the plane perpendicular to the $\overrightarrow{\Omega}$ axis



• $\overrightarrow{\Omega}$ is directed out of the center of the cirle and $|\delta \hat{k}'| = \sin \alpha \delta \theta$.

$$\left| \frac{D_a \hat{k}'}{Dt} \right| = \lim_{\delta t \to 0} \frac{\left| \delta \hat{k}' \right|}{\delta t} = \lim_{\delta t \to 0} \sin \alpha \frac{\delta \theta}{\delta t} = \Omega \sin \alpha$$

But
$$\left| \overrightarrow{\Omega} \times \hat{k}' \right| = \left| \overrightarrow{\Omega} \right| \left| \hat{k}' \right| \sin \alpha$$
, so $\left| \frac{D_a \hat{k}'}{Dt} \right| = \left| \overrightarrow{\Omega} \times \hat{k}' \right|$



- $D_a \hat{k}'/Dt$ has direction of $\delta \hat{k}'$ as $\delta t \to 0$
- $\delta \hat{k}'$ has direction of $\overrightarrow{\Omega} \times \hat{k}'$ ($\delta \hat{k}'$ is perpendicular to $\overrightarrow{\Omega}$ and perpendicular to \hat{k}' and points in the direction of increasing θ)
- Thus, $D_a \hat{k}'/Dt$ has direction $\overrightarrow{\Omega} \times \hat{k}'$
- Since $D_a \hat{k}'/Dt$ has the same magnitude and direction as $\overrightarrow{\Omega} \times \hat{k}'$, it must be $\overrightarrow{\Omega} \times \hat{k}'$
- Thus

$$\frac{D_a \hat{k}'}{Dt} = \vec{\Omega} \times \hat{k}'$$

and similarly,

$$\frac{D_a \hat{i}'}{Dt} = \vec{\Omega} \times \hat{i}' \qquad \text{and} \qquad \frac{D_a \hat{j}'}{Dt} = \vec{\Omega} \times \hat{j}'$$



• We can now write our expression as:

$$\frac{D_{a}\overrightarrow{A}}{Dt} = \frac{D\overrightarrow{A}}{Dt} + A'_{x}\overrightarrow{\Omega} \times \hat{i}' + A'_{y}\overrightarrow{\Omega} \times \hat{j}' + A'_{z}\overrightarrow{\Omega} \times \hat{k}'$$

$$\frac{D_{a}\overrightarrow{A}}{Dt} = \frac{D\overrightarrow{A}}{Dt} + \overrightarrow{\Omega} \times \left(A'_{x}\hat{i}' + A'_{y}\hat{j}' + A'_{z}\hat{k}'\right)$$

$$\frac{D_{a}\overrightarrow{A}}{Dt} = \frac{D\overrightarrow{A}}{Dt} + \overrightarrow{\Omega} \times \overrightarrow{A}$$

 In words, this says that the total derivative as observed in a non-rotating reference frame is equal to the total derivative as observed in a rotating reference frame plus the apparent force arising from rotation.

- Let's consider the vector form of $\vec{F} = m\vec{a}$, or $\vec{a} = \vec{F}/m$, which holds for an inertial reference frame.
- \vec{a} is the acceleration observed in an inertial reference frame:

$$\vec{a} = \frac{D_a \vec{U}_a}{Dt}$$

where $\overrightarrow{U}_a = D_a \overrightarrow{r}/Dt$ is the absolute velocity of an air parcel and \overrightarrow{r} is the position vector.

• Now we apply the previous relationship and set $\overrightarrow{A} = \overrightarrow{r}$

$$\frac{D_a \vec{r}}{Dt} = \frac{D \vec{r}}{Dt} + \vec{\Omega} \times \vec{r}$$

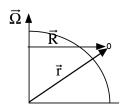
$$\vec{U}_a = \underbrace{\vec{U}}_{l} + \underbrace{\vec{\Omega} \times \vec{r}}_{l}$$

where (I) is the relative velocity of the air parcel and (II) is the velocity of solid-body rotation at location \vec{r} .so



• Now let's set $\vec{A} = \vec{U}_a$:

$$\begin{split} \frac{D_a \overrightarrow{U}_a}{Dt} &= \frac{D \overrightarrow{U}_a}{Dt} + \overrightarrow{\Omega} \times \overrightarrow{U}_a \\ &= \frac{D}{Dt} \left(\overrightarrow{U} + \overrightarrow{\Omega} \times \overrightarrow{r} \right) + \overrightarrow{\Omega} \times \left(\overrightarrow{\Omega} \times \overrightarrow{r} \right) \\ &= \frac{D \overrightarrow{U}}{Dt} + \overrightarrow{\Omega} \times \underbrace{\frac{D \overrightarrow{r}}{U}}_{\overrightarrow{U}} + \overrightarrow{\Omega} \times \overrightarrow{U} + \underbrace{\overrightarrow{\Omega} \times \left(\overrightarrow{\Omega} \times \overrightarrow{r} \right)}_{-\Omega^2 \overrightarrow{R}} \end{split}$$





• The final form:

$$\underbrace{\frac{D_a \vec{U}_a}{Dt}}_{1} = \underbrace{\frac{D\vec{U}}{Dt}}_{2} + \underbrace{2\vec{\Omega} \times \vec{U}}_{3} - \underbrace{\Omega^2 \vec{R}}_{4}$$

- 1 acceleration in an inertial reference frame
- 2 acceleration in a non-inertial reference frame
- Coriolis acceleration
- 4 centripetal acceleration

Plus this into Newton's 2nd Law:

$$\frac{D_a \vec{U}_a}{Dt} = \frac{\vec{F}}{m}$$

$$\frac{D\vec{U}}{Dt} + 2\vec{\Omega} \times \vec{U} - \Omega^2 \vec{R} = \frac{\vec{F}}{m}$$

- Recall that \vec{F}/m is the sum of all the fundamental forces per unit mass
 - PGF: $-\frac{1}{\rho} \vec{\nabla} p$ friction: $\nu \vec{\nabla}^2 \vec{U}$

 - gravitational: $\overrightarrow{a^*}$



Newton's 2nd Law becomes:

$$\frac{D\vec{U}}{Dt} + 2\vec{\Omega} \times \vec{U} - \Omega^2 \vec{R} = -\frac{1}{\rho} \vec{\nabla} p + \vec{g^*} + \nu \vec{\nabla}^2 \vec{U}$$
$$\frac{D\vec{U}}{Dt} = -\frac{1}{\rho} \vec{\nabla} p - 2\vec{\Omega} \times \vec{U} + \underbrace{\vec{g^*} + \Omega^2 \vec{R}}_{\vec{g}} + \nu \vec{\nabla}^2 \vec{U}$$

We arrive at the vector equation of motion in a rotating reference frame:

$$\boxed{ \frac{D \overrightarrow{U}}{D t} = -\frac{1}{\rho} \overrightarrow{\nabla} p - 2 \overrightarrow{\Omega} \times \overrightarrow{U} + \overrightarrow{g} + \nu \overrightarrow{\nabla}^2 \overrightarrow{U} }$$



$$\boxed{ \underbrace{\frac{D\vec{U}}{Dt}}_{1} = \underbrace{-\frac{1}{\rho}\vec{\nabla}p - 2\vec{\Omega} \times \vec{U}}_{2} + \underbrace{\vec{g}}_{4} + \nu\vec{\nabla}^{2}\vec{U}}_{5} }$$

- 1 acceleration in a rotating reference frame
- pressure gradient force
- 6 Coriolis force
- gravity
- 6 friction

